

AXIAL VIBRATION OF A CANTILEVER BEAM WITH A TIP MASS

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ABSTRACT

In this study free axial vibration of a cantilever beam with a tip mass is analyzed. The boundary conditions are written for the fixed end and the end with the tip mass after the differential equation of motion is solved by separation of variables method. The frequency values for the first three vibration modes of the beam are obtained for various values of concentrated mass and presented in the tables. The frequency values for the concentrated mass values of zero are also compared with the ones of cantilever beam without tip mass and nearly the exact values are obtained with negligible error percentages.

Keywords: Axial vibration, Cantilever Beam, Tip Mass, Frequency

1. Introduction

In practice, the representation of a beam by a discrete model is an idealized model; however, in fact, beams have continuously distributed mass and elasticity. Mostly, especially for the axially vibration, beams are modeled as continuous systems having infinite number of degreed of freedom [1-10].

In this study, the free vibration analysis of a uniform axially vibrating cantilever beam with a tip mass is made. The differential equation of motion of the axially vibrating beam is solved by separation of variables method [11] and the displacement function is obtained. The boundary conditions are written for the fixed end and the tip mass. The natural frequencies for the first three modes are obtained for the various values of the tip mass. The results obtained for the tip mass value of zero are compared with the frequency values of the cantilever beam without a tip mass. The axially vibrating beam considered in the study is assumed to be homogeneous and isotropic.

2. Solution of Equation of Motion for an Axially Vibrating Beam

An axially vibrating beam, given in Figure.1, with the distributed mass m, the length L, the modulus of elasticity E, the cross-section area A and the axial rigidity AE has a dimensional differential equation of motion for free vibration as [12]

$$\xrightarrow{u(x,t)} x$$
m, L, AE

Figure 1: An Axially Vibrating Beam

$$\frac{\partial^2 u(x,t)}{\partial x^2} - \frac{m}{AE} \frac{\partial^2 u(x,t)}{\partial t^2} = 0$$

(1)

where u(x,t) is the displacement function of the beam in terms of both displacement x and time t. Application of the separation of variables method to Eq. (1) as in the form of Eq. (2) is commonly used in vibration analysis of beams.

$$u(x,t) = X(x) \cdot T(t) = X(x) \cdot [A \cdot \sin(\omega t) + B \cdot \cos(\omega t)]$$
(2)

In Eq. (2), X(x) is the eigenfunction named as shape function, T(t) is time function, ω is the eigenvalue of the solution named as natural frequency and A, B are constants.

The derivatives used in Eq. (1) can, therefore, be written as

$$\frac{\partial^2 u(x,t)}{\partial x^2} = u''(x,t) = X''(x). [A.\sin(\omega t) + B.\cos(\omega t)] = X''(x).T(t)$$
(3)

$$\frac{\partial^2 u(x,t)}{\partial t^2} = \ddot{u}(x,t) = X(x). (-\omega^2)[A.\sin(\omega t) + B.\cos(\omega t)] = -\omega^2.X(x).T(t)$$

where $\binom{n}{2}$ and $(\ddot{})$ denote the second order derivative due to x and t, respectively. Substitution of Eq. (3) and Eq. (4) in Eq. (1) gives the governing equation of motion in the form as

$$X^{''}(x).T(t) + \frac{m\omega^2}{AE}X(x).T(t) = 0$$
 $X^{''}(x) + \frac{m\omega^2}{AE}X(x) = 0$ $0 \le x \le L$

for
$$\alpha^2 = \frac{m\omega^2}{AE}$$
 $X''(x) + \alpha^2 X(x) = 0$

(6)

(5)

(4)

$$D^2 + \alpha^2 = 0 \quad \to \quad D_{1,2} = \pm i\alpha \tag{7}$$

$$X(x) = C_1 \sin(\alpha x) + C_2 \cos(\alpha x)$$
(8)

Eq. (8) gives the shape function of the axially vibrating beam due to the displacement variable, x. Therefore, from Eq. (2), the displacement function of the axially vibrating beam has the form of Eq. (9).

$$u(x,t) = X(x).T(t) = [C_1.sin(\alpha x) + C_2.cos(\alpha x)].T(t)$$
(9)

3. Boundary Conditions



Figure 2: Axially Vibrating a Cantilever Beam with a Tip Mass

Two boundary conditions have to be written for the cantilever beam with a tip mass in Figure.2 since two integration constants (C_1 , C_2) are obtained in the solution of second order differential equation of motion. The boundary conditions written for the left and the right ends of axially vibrating beam are given, respectively, as [13]

for x=0
$$u(x = 0, t) = 0$$

(10)
for x=L $N(x = L, t) = AEu'(x = L, t) = -M.\ddot{u}(x = L, t)$
(11)

where M is the tip mass value and N(x,t) is the axial force. If Eq. (9) and its derivative are substituted into Eq. (10) and Eq. (11) one gets the following relation between the coefficient matrix and the integration constants.

$$[k]. \{C\} = 0; \quad [\propto .\cos(\propto L) - \alpha_M^2 . \sin(\propto L)]. \{C_1\} = \{0\}$$

$$(12)$$

$$|k| = |\propto .\cos(\propto L) - \alpha_M^2 . \sin(\propto L)| = 0$$

$$(13)$$

where $\propto_M^2 = \frac{M \cdot \omega^2}{AE}$. For non-trivial solution equating the determinant of the coefficient matrix of Eq. (12) to zero, as in Eq. (13), will give the eigenfrequencies of the axially vibrating cantilever beam with a tip mass. These frequencies are computed by a program written by the author considering the secant method [14].

4. Numerical Analysis

The first three natural frequencies of the axially vibrating cantilever beam with a tip mass are calculated for M values of 0, 10^{-10} , 10^{-9} , 10^{-8} , 10^{-6} , 10^{-4} , 10^{-2} , 10^{-1} , 10^{0} , 10^{1} and 10^{2} , the beam length of L=1 m. and the modulus of elasticity of E=2100000 kg/cm². IPB-100, IPB-300 and IPB-600 profiles are used for numerical analysis with the mechanical properties given in Table 1 where h is height, G is weight per length, A is cross-section area and AE is axial rigidity of the corresponding profile. The distributed mass of the beam m is calculated from G/g as g being the acceleration of gravity with the value of 981 cm/sn².

| Profile | h | G | Α | AE | |
|---------|------|---------|----------------------------|-----------|--|
| | (cm) | (kg/cm) | (cm ²) | (kg) | |
| IPB100 | 10 | 0.081 | 10.3 | 21630000 | |
| IPB300 | 30 | 0.422 | 53.8 | 112980000 | |
| IPB600 | 60 | 1.22 | 156 | 327600000 | |

Table 1: The Mechanical Properties of the Profiles Used in This Study

The frequency values computed due to different values of the tip mass are presented in Table 2, Table 3 and Table 4 for, respectively, IPB-100, IPB-300 and IPB-600. The frequency values obtained for M=0 are compared with the exact frequency values obtained from the frequency equation of cantilever beam without a tip mass in the first row.

| | | ω1 | ω_2 | ω ₃ |
|---------------------------------------|---|-----------|------------|----------------|
| Exact Frequencies for Cantilever Beam | $\omega_i = \frac{(2n_i - 1)}{2} \frac{\pi}{L} \sqrt{\frac{AE}{m}}$ | 8039.7053 | 24119.1160 | 40198.5266 |
| | 0 | 8039.7053 | 24119.1160 | 40198.5266 |
| | 10⁻¹⁰ | 8039.7052 | 24119.1157 | 40198.5261 |
| | 10 ⁻⁹ | 8039.7044 | 24119.1131 | 40198.5218 |
| | 10⁻⁸ | 8039.6956 | 24119.0868 | 40198.4779 |
| | 10 ⁻⁶ | 8039 | 24117 | 40194 |
| Μ | 10 ⁻⁴ | 7944 | 23831 | 39719 |
| | 10 ⁻² | 4099 | 17306 | 32815 |
| | 10 ⁻¹ | 1451 | 16213 | 32226 |
| | 10 ⁰ | 465 | 16093 | 32166 |
| | 10 ¹ | 148 | 16081 | 32160 |
| | 10 ² | 47 | 16080 | 32159 |

Table 2: Frequencies Computed for Different Values of the Tip Mass for IPB-100

Table 3: Frequencies Computed for Different Values of the Tip Mass for IPB-300

| | | ω ₁ | ω ₂ | ω ₃ |
|---------------------------------------|---|----------------|----------------|----------------|
| Exact Frequencies for Cantilever Beam | $\omega_i = \frac{(2n_i - 1)}{2} \frac{\pi}{L} \sqrt{\frac{AE}{m}}$ | 8050.0567 | 24150.1702 | 40250.2836 |
| М | 0 | 8050.0567 | 24150.1702 | 40250.2836 |
| | 10⁻¹⁰ | 8050.0567 | 24150.1701 | 40250.2835 |
| | 10⁻⁹ | 8050.0565 | 24150.1696 | 40250.2827 |
| | 10 ⁻⁸ | 8050.0549 | 24150.1646 | 40250.2743 |
| | 10 ⁻⁶ | 8050 | 24150 | 40250 |
| | 10⁻⁴ | 8032 | 24095 | 40157 |
| | 10 ⁻² | 6567 | 20333 | 35076 |
| | 10 ⁻¹ | 3139 | 16770 | 32547 |

| 10 ⁰ | 1056 | 16170 | 32236 |
|----------------------------|------|-------|-------|
| 10¹ | 336 | 16108 | 32204 |
| 10² | 107 | 16770 | 32201 |

| | | ω1 | ω ₂ | ω3 |
|--|---|-----------|----------------|------------|
| Exact Frequencies for Cantilever Beam | $\omega_i = \frac{(2n_i - 1)}{2} \frac{\pi}{L} \sqrt{\frac{AE}{m}}$ | 8062.0672 | 24186.2015 | 40310.3358 |
| | 0 | 8062.0672 | 24186.2015 | 40310.3358 |
| | 10⁻¹⁰ | 8062.0672 | 24186.2015 | 40310.3358 |
| | 10⁻⁹ | 8062.0671 | 24186.2013 | 40310.3355 |
| | 10⁻⁸ | 8062.0665 | 24186.1996 | 40310.3329 |
| | 10 ⁻⁶ | 8062 | 24186 | 40310 |
| Μ | 10 ⁻⁴ | 8056 | 24167 | 40278 |
| | 10⁻² | 7465 | 22451 | 37580 |
| | 10⁻¹ | 4769 | 17884 | 33223 |
| | 10 ⁰ | 1774 | 16325 | 32350 |
| | 10 ¹ | 572 | 16145 | 32259 |
| | 10² | 181 | 16127 | 32250 |

Table 4: Frequencies Computed for Different Values of the Tip Mass for IPB-600

5. Conclusions

In this study free axial vibration of a cantilever beam with a tip mass is made. The natural frequency values are obtained for different values of the tip mass and presented in tables. It can be seen from Tables 2, 3 and 4 that as the tip mass values increase from zero through a value of 10^{-6} for all profiles considered in this study the frequency values gently decrease, however, after the mentioned values the frequency values begin to decrease dramatically and rapidly. Increasing the height of the beam section causes an increase in frequency values.

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